# The Breakthrough Listen Search for Intelligent Life: Observations of 1327 Nearby Stars Over 1.10–3.45 GHz

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# Abstract

Breakthrough Listen (BL) is a 10 year initiative to search for signatures of technologically capable life beyond Earth via radio and optical observations of the local universe. A core part of the BL program is a comprehensive survey of 1702 nearby stars at radio wavelengths (1-10 GHz). Here, we report on observations with the 64 m CSIRO Parkes radio telescope in New South Wales, Australia, and the 100 m Robert C. Byrd Green Bank radio telescope in West Virginia, USA. Over 2016 January to 2019 March, a sample of 1138 stars was observed at Green Bank using the 1.10–1.90 GHz and 1.80–2.80 GHz receivers, and 189 stars were observed with Parkes over 2.60–3.45 GHz. We searched these data for the presence of engineered signals with Doppler-acceleration drift rates between  $\pm 4 \,\mathrm{Hz}\,\mathrm{s}^{-1}$ . Here, we detail our data analysis techniques and provide examples of detected events. After excluding events with characteristics consistent with terrestrial radio interference, we are left with zero candidates. That is, we find no evidence of putative radio transmitters above  $2.1 \times 10^{12}$  W, and  $9.1 \times 10^{12}$  W for Green Bank and Parkes observations, respectively. These observations constitute the most comprehensive search over 1.10–3.45 GHz for technosignatures to date. All data products, totaling  $\sim$ 219 TB, are available for download as part of the first BL data release (DR1), as described in a companion paper.

Unified Astronomy Thesaurus concepts: Astrobiology (74); Technosignatures (2128); Search for extraterrestrial intelligence (2127)

# 1. Introduction

If we are to detect life beyond Earth in the next few decades, it will be by one of three ongoing efforts. We may find life in the solar system by physically examining the environment of our planetary neighbors and their moons. Optical spectroscopy may detect biosignatures in the atmospheres of nearby exoplanets, indicating the presence of life. Or, we may detect evidence of advanced life via technosignatures: signals of engineering that are discernible from astrophysical processes.

These three methods are complementary, as they probe different manifestations of life at different distance scales and timescales in life evolution. The last approach is known as the Search for Extraterrestrial Intelligence (SETI) and is the only method that can conceivably detect signals from beyond the nearest stars with current or near-term technology. SETI seeks not just to detect signs of life, but also to constrain the probability of the emergence of intelligence life: whether we are the sole inhabitants of the universe, or whether it is ours to share.

Radio searches for extraterrestrial intelligence (ETI) have been ongoing since the 1960s (Drake 1961). The sensitivity and speed of SETI searches are intimately tied to our own technological capability; as technology progresses, so too do the capabilities and

sensitivity of radio telescopes. Of particular importance in this regard are the digital instruments used in radio SETI searches. The instantaneous bandwidth of these systems has expanded from hundreds of hertz (Drake 1961), to kilohertz (e.g., Werthimer et al. 1985; Horowitz & Sagan 1993), to megahertz (Werthimer et al. 1995; Tarter 1996), and through to tens of gigahertz (MacMahon et al. 2018)—a factor of  $10^8$ —over the course of roughly 40 yr. The search has also expanded from single-dish radio telescopes to interferometers (e.g., Welch et al. 2009; Rampadarath et al. 2012; Tingay et al. 2016, 2018a, 2018b; Gray & Mooley 2017), optical (Wright et al. 2001; Reines & Marcy 2002; Howard et al. 2004; Stone et al. 2005; Howard et al. 2007; Tellis & Marcy 2015) and infrared wavelengths (Slysh 1985; Carrigan 2009; Wright et al. 2014a, 2014b; Griffith et al. 2015). In tandem with increased frequency coverage, the sensitivity and field of view of telescopes continue to increase, allowing ever more exquisite measurements to be made.

# 1.1. Breakthrough Listen

The Breakthrough Listen (BL) initiative represents the current state-of-the-art for SETI search strategies and capabilities. BL is a 10 year initiative to search for technosignatures at



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 Table 1

 Details of the Receivers Used Here

Telescope	Receiver	Frequency (GHz)	T <sub>sys</sub> (K)	SEFD (Jy)
Green Bank	L band	1.10-1.90	25	10
Green Bank	S band	1.80-2.80	25	10
Parkes	10 cm	2.60-3.45	35	34

radio and optical wavelengths (Worden et al. 2017). Having commenced observations in 2016, the program expands the capability of existing telescopes for SETI observations by installing wide-bandwidth data recording and analysis systems capable of recording raw digitizer samples direct to disk (see MacMahon et al. 2018; Price et al. 2018; Lebofsky et al. 2019). In its initial years, BL is conducting observations with the 100 m Robert C. Byrd Green Bank (henceforth GBT) radio telescope in West Virginia, USA, and the 64 m CSIRO Parkes radio telescope in New South Wales, Australia. New digital systems have been installed at both telescopes to allow for capture of voltage data across the full bandwidth provided by the receivers of the two telescopes (MacMahon et al. 2018; Price et al. 2018). At optical wavelengths, the Automated Planet Finder telescope in California, USA (Radovan et al. 2014), is conducting a search for laser lines in high-resolution spectra (e.g., Lipman et al. 2019). BL is currently conducting a survey of several thousand nearby stars, 100 nearby galaxies, and the Galactic plane; further details may be found in Isaacson et al. (2017). An initial analysis of 692 stars, over 1.1-1.9 GHz, is presented in Enriquez et al. (2017); no high-duty-cycle narrowband radio transmissions with equivalent isotropic radiated power (EIRP) of  $>10^{13}$  W were found in this sample. Here, we present an analysis of 1327 star targets taken over

Here, we present an analysis of 1327 star targets taken over 1.10-3.45 GHz (*L* band and *S* band), including reanalysis of the observations of 692 stars detailed in Enriquez et al. (2017). In addition to covering greater bandwidth, we improve on the Enriquez et al. (2017) limits by using a lower signal-to-noise ratio (S/N) cutoff (10 versus 25), a larger range of drift rates ( $\pm 4$  Hz s<sup>-1</sup> versus  $\pm 2$  Hz s<sup>-1</sup>), and enhanced signal identification logic.

This paper is organized as follows. In Section 2, a summary of observations is given. Data analysis strategies are detailed in Section 3. Results are given in Section 4, followed by a discussion and conclusions. An explanation of the data formats and archiving strategy is given in a companion paper (Lebofsky et al. 2019).

# 2. Observations

We used the Green Bank and Parkes telescopes to observe nearby stars at frequencies between 1.10 and 3.45 GHz (see Table 1). This section provides a summary of the stellar targets observed, details of the two telescopes used, observational details, and an overview of data products. A full listing of observed stars can be found at seti.berkeley.edu/listen2019.

# 2.1. Star Sample

We observed nearby stellar targets chosen from the Isaacson et al. (2017) 1702 star sample, with the Green Bank and Parkes radio telescopes. The Isaacson et al. (2017) sample is comprised of stars selected from the RECONS and Gliese catalogs of nearby stars (Gliese & Jahreiss 1995) and the well-characterized *Hipparcos* catalog (Perryman et al. 1997). The sample is constructed to contain all stars within 5 pc in the Gliese and RECONS catalogs and a broad sampling of main-sequence stars within 5–50 pc from *Hipparcos*.

For observations with Parkes, the Isaacson et al. (2017) sample was augmented with a small number of recently discovered brown dwarfs and other stars within 5 pc, below a decl. of  $-15^{\circ}$ ; these are detailed in Table 2. In total, 1327 distinct primary "A" star targets were observed. Their distribution on the sky, and a histogram of distances, are shown in Figures 1 and 2, respectively. It should be noted that this is still not a complete list, and the continuing discovery of nearby low-mass stars necessitates periodic revisiting of our volume-limited sample. Also, we note that not all stars were observed with each single receiver. The number of stars per receiver is broken down below and summarized in Table 3.

# 2.2. Observing Strategy

At both GBT and Parkes, we employ an observing strategy whereby a target is observed for five minutes ("ON" source), then a reference location is observed for five minutes ("OFF" source). This ON-OFF strategy is repeated three times for each target, taking a total of 30 minutes (plus slew time). This strategy is used to allow for discrimination of bona fide signals of interest from radio interference (RFI): any signal that appears in both ON and OFF pointings at power levels inconsistent with the known off-axis gain of the telescope is considered RFI. To further discriminate RFI-induced false positives, we enforce that signals must appear in all three ON pointings, or in other words, that the signals are continuous throughout the observation. Further discrimination may then be done by enforcing that signals exhibit a nonzero Doppler acceleration (see Section 3) and by cross-reference between observations.

Following Enriquez et al. (2017), we refer to a strategy of applying a constant offset for OFF positions as ABABAB; the strategy of interspersing three different nearby "secondary" targets is referred to as ABACAD. At Parkes, an ABABAB strategy is used with a fixed 0°.5 decl. offset (~3 FWHM beamwidths at the lowest observing frequency); at GBT, an ABACAD strategy is predominantly used. For GBT ABACAD observations, nearby stars selected from the Hipparcos catalog are used for OFF-source pointings (these "secondary" stars may be searched for technosignature or flare activity at a later date). The Hipparcos stars are chosen to be between 1.2 and 3.6 degrees away from the primary target (>8 FWHM beamwidths at the lowest observing frequency). This separation was chosen to be sufficiently far from the primary beam, within a reasonable slew time and encompassing an area that likely holds three Hipparcos stars. A search of the secondary targets is also possible, but outside the scope of this paper.

# 2.3. Green Bank Telescope

The GBT is a 100 m radio telescope located in West Virginia, USA (38°25′59″236 N,79°50′23″406 W), operated by the Green Bank Observatory. The telescope is located within a federally protected "Radio Quiet" zone, in which most radio transmissions are prohibited, to minimize radio frequency interference (RFI). Approximately 20% of the annual observing time for the GBT is dedicated to BL. The GBT has an operational range of 0.3–110 GHz, depending on the receiver

 Table 2

 Stars within 5 pc Added to the Isaacson et al. (2017) Sample for Improved Volume Completeness Below  $-15^{\circ}$  Decl.

Star	Epoch	R.A.	Decl.	Distance (pc)	Spectral Type	Reference
DENIS J025503.3-470049	2000	2:55:03.6	-47:00:51.0	4.9	L8/L9	Cutri et al. (2003)
SCR J1845-6357A	2000	18:45:05.3	-63:57:47.1	3.9	M8.5	Biller et al. (2006)
WISE J035000.32-565830.2	2000	3:50:00.3	-56:58:30.2	5.4	Y1	Cutri et al. (2012)
WISE J053516.80-750024.9	2000	5:35:16.8	-75:00:24.9	2.0	Y1	Cutri et al. (2012)
WISE J104915.57-531906.1	2000	10:49:18.9	-53:19:10.1	2.0	L7.5, T0.5	Cutri et al. (2003)
WISEA J154045.67-510139.3	2000	15:40:45.7	-51:01:39.3	4.6	M7	Kirkpatrick et al. (2014)

 Table 3

 Summary of GBT S, GBT L, and Parkes 10 cm Observations

Receiver	No. Cadences	No. Targets	Hr
GBT L	1013	882	506.5
GBT S	1076	1005	538.0
Parkes 10 cm	966	189	483.0

equipped during observation. For the analyses detailed here, we used the 1.10-1.90 GHz (*L*-band) receiver and the 1.80-2.80 GHz (*S*-band) receiver, both with a system temperature of 20 K, resulting in a system-equivalent flux density (SEFD) of 10 Jy. The *L* band contains a user-selectable notch filter band between 1.20 and 1.34 GHz (which is always used in BL observations), and the *S* band contains a permanently installed superconducting notch filter band between 2.3 and 2.36 GHz (NRAO 2019).

At the time of writing, the nearby star observation program at GBT is currently focusing on observations with the 4.0–7.8 GHz (*C*-band) and 7.8–12.3 GHz (*X*-band) receivers. Completion of these programs are at about 80% and 60%, respectively. Use of the 18–27.5 GHz (KFPA-band) receiver for BL observations is being commissioned, with the full 10 GHz bandwidth of raw voltage data produced by this receiver available to the current BL backend (MacMahon et al. 2018). Analysis of these data is expected to be included in future publications.

A total of 12,504 five-minute observations with the GBT ( $\sim$ 1044 hr) are used for this work, conducted over the period 2016 January 1–2019 March 23 (MJD 57388–58565), summarized in Table 3. Out of these observations, 6042 were carried out with the *L*-band receiver (1013 cadences) and 6456 were carried out with the *S*-band receiver (1076 cadences). Due to a small number of repeated observations, where the star was selected more than once from the target database, a total of 1138 primary "A" stars were observed: 749 at both bands, 882 at the *L* band, and 1005 at the *S* band.

### 2.4. Parkes Radio Telescope

The CSIRO Parkes radio telescope is a 64 m telescope located in New South Wales, Australia (32°59′59″8S, 148°15′44″3E). As with the GBT, Parkes is equipped with a suite of receivers, which covers 0.6–26.0 GHz. Over the period 2016 October– 2021 September, a quarter of the annual observation time of the Parkes 64 m radio telescope has been dedicated to the BL program. The analyses detailed here are from data taken with the 10 cm component of the Parkes "10–50" receiver, which covers 2.60–3.45 GHz. This receiver has a nominal system temperature of 35 K, with a corresponding SEFD of 34 Jy.

In contrast to the GBT, the Parkes observations of the nearby star sample used herein include multiple epochs. This was motivated primarily by technical concerns: first, the Parkes system was deployed in stages over a period of several months when the receiver availability varied (two receivers can be installed in the focus cabin at Parkes, and the choice of receivers is motivated both by technical availability and scheduling constraints); second, for an isolated transmitter on the surface of a rotating body, we would expect intermittent behavior, which presents a potential opportunity for increasing the probability of interception with repeated observation; and finally, use of the 0.7–4.0 GHz Ultra-Wideband Low ("UWL") receiver (Hobbs et al. 2020), is planned for future observations, which is the appropriate complement to observations with GBT.

At Parkes, a total of 483.0 hr of observations are used for this work, over the period 2016 November 16–2018 January 19 (MJD 57708–58137). During this time, 966 cadences covering a total of 189 targets were observed: 183 from the 1702 star sample and 6 additional nearby stars (Table 2).

Work with other receivers, including the 21 cm multibeam (MB; Staveley-Smith et al. 1996), in addition to the already mentioned UWL (Manchester 2015), is ongoing and not included here.

### 2.5. Data Reduction Pipeline

A comprehensive overview of the BL data products and reduction pipelines is given by Lebofsky et al. (2019); here we provide a brief summary. Both GBT and Parkes use the same hardware and firmware to sample the incoming analog signals from the receiver, which we refer to as the signal processing "front end." This hardware, a 5 Gsample/s, 8 bit digitizer and field-programmable gate array (FPGA) processing board, is provided by the Collaboration for Astronomy Signal Processing and Electronics Research (CASPER; Hickish et al. 2016). Detailed instrument descriptions are provided by MacMahon et al. (2018) and Price et al. (2018).

During observations, the front end samples the dual-polarization receiver output at 8 bit resolution, then applies a polyphase filterbank to coarsely channelize the data into  $\sim 2.92$  MHz bands, running firmware detailed by Prestage et al. (2015). The front-end FPGA boards output channelized data over 10 Gb Ethernet to a cluster of high-performance compute nodes, each of which captures 187.5 MHz of dual-polarization data. The compute nodes write 8 bit voltage-level products to disk in raw format.<sup>13</sup> Each node is equipped with a Graphics Processing Unit (GPU), which is used to convert the voltage-level raw files into spectral data products, stored in filterbank format.

A total of three filterbank data products are generated: a high-spectral-resolution product with frequency and time resolution of  $\sim 2.79$  Hz and  $\sim 18.25$  s respectively, a midresolution product ( $\sim 2.86$  kHz,  $\sim 1.07$  s), and a high-time-resolution

<sup>&</sup>lt;sup>13</sup>The raw, filterbank, and hdf5 formats are detailed by Lebofsky et al. (2019).

product ( $\sim$ 366 kHz,  $\sim$ 349  $\mu$ s). Here, as we are searching for the presence of narrowband signals, we analyze only the high-spectral-resolution product.

After observations are completed, the spectral products from each compute node are combined into a single frequencycontiguous file and converted into hdf5 format. All data analyzed here are available online<sup>14</sup>; final data volumes are 142 TB and 77 TB for Green Bank and Parkes, respectively.

# 3. Methods

For a receiver fixed on Earth, any transmitter not also on Earth's surface (or in geosynchronous orbit) will exhibit a timedependent Doppler shift due to relative acceleration between the transmitter (ET) and receiver. The maximum Doppler shift in frequency,  $\Delta \nu_{max}$ , depends upon the relative velocity  $\Delta v$  of the transmitter and transmitted frequency  $\nu_{\text{ET}}$ :

$$\Delta \nu_{\max} = \nu_{\rm ET} \left( \frac{\Delta \nu}{c} \right). \tag{1}$$

Over short ( $\sim$ 5 minute) durations, the change in frequency is well approximated as a linear function,

$$\nu(t) = \nu_0 + \dot{\nu}t,\tag{2}$$

where  $\nu_0$  is the frequency at t = 0 and  $\dot{\nu}$  is the shift in frequency (units of  $Hz s^{-1}$ ) due to Doppler motion, or drift rate. If after a time  $t_{obs}$  the product  $\dot{\nu} \times t_{obs}$  is greater than the channel bandwidth B, signal power will be "smeared" across  $N = \dot{\nu} \times t_{\rm obs}/B$  channels, lowering the detected signal to noise (S/N) in each channel by a factor  $N^{1/2}$ . This effect can be compensated for, if the observation is split into subintegrations, by applying a shift to each subintegration before integrating to form a final spectrum. Usually, the subintegration smearing is not corrected for-this approach is known as an incoherent search. If the drift rate is not known, a search across a range of trial drift rates can be conducted to identify the drift rate that optimizes detection; this search can be done by a brute-force approach or other means. Our detection algorithm turboSETI uses a tree search algorithm which is optimized computationally (Enriquez et al. 2017 and references therein).

The maximum Doppler drift due to a body's rotation is given by

$$\dot{\nu} = \frac{4\pi^2 R}{P^2} \frac{\nu_0}{c},$$
(3)

where *c* is the speed of light, *P* is rotational (or orbital) period, and *R* is the body (or orbit) radius. At the lower and upper frequency limits of our data set (1.1-3.4 GHz), Earth's daily rotation corresponds to drift rates of magnitude  $0.12-0.38 \text{ Hz s}^{-1}$ ; Earth's 1 au orbit imparts  $0.02-0.06 \text{ Hz s}^{-1}$ . Here, we search drift rates between -4 and  $+4 \text{ Hz s}^{-1}$ , corresponding to fractional drift rates of 3.64 nHz at 1.1 GHz, and 1.15 nHz at 3.45 GHz. This rate allows for a wide range of planetary radii, spin periods, and orbital periods: Oliver & Billingham (1971) advocate a 1 nHz rate, based on that expected from an Earth-like planet with an 8 hr day. More recently, in simulations of planetary formation models, Miguel & Brunini (2010) find that a majority of planetary primordial rotation periods for  $<10M_{\oplus}$  planets fall between 10 and 10,000 hr; this sets a 0.65 nHz maximum rate. Note that a putative transmitter

that is not gravitationally bound would still exhibit a Doppler drift imparted by Earth's rotation.

### 3.1. Dedoppler Search Analysis

Following Enriquez et al. (2017), we use the *turbo*SETI incoherent dedoppler code (Enriquez & Price 2019)<sup>15</sup> to search our high-resolution ( $\sim$ 2.79 Hz) Stokes-*I* data for drifting narrowband signals.

*turbo*SETI applies an efficient tree search algorithm, based on Taylor (1974) and Siemion et al. (2013). The tree algorithm efficiently computes the integrals over straight line paths. It is similar to the Hough transform (e.g., Leavers 1992), which also computes integrals over straight line paths, but also applies an edge-detection step to convert input data into a binary image. The Hough transform itself can be thought of as a two-dimensional discrete Radon transform (Götz & Druckmüller 1995). The tree dedispersion algorithm accelerates searches by reuse of redundant computations involved when searching similar slopes, which reduces the number of additions required to  $n \log_2 n$  from  $n^2$ , where *n* is equal to both the number of spectra and number of slopes searched.

By the radiometer equation, the noise within a coarse channel (without RFI present) follows a chi-squared distribution, as the digitized voltages are well approximated as zeromean Gaussian random processes (Thompson et al. 2017). Here, we compute root mean square noise levels from the 90th central percentile of the power values to mitigate outliers in the distribution, due to the presence of narrowband features and rolloff imparted by the shape of the polyphase filter.

The number of discrete frequency drift rates that are searched by *turbo*SETI depends on  $\dot{\nu}_{min} = B/t_{obs}$ , roughly 0.01 Hz s<sup>-1</sup> for our high-frequency resolution data products. For our search to rates ±4 Hz s<sup>-1</sup>, roughly 800 Doppler trials are performed.

We ran *turbo*SETI on all files, searching drift rates  $\pm 4$  Hz, for narrowband signals with an S/N  $\ge$  10. We parallelized processing tasks across nodes using a code called TIDDALIK,<sup>16</sup> which distributes and executes tasks across nodes. As *turbo*SETI runs on a single file, processing is "pleasingly parallel" and can be run without intertask communication.

*turbo*SETI produces a list of "hits," i.e., detections above a given S/N, in a .dat plain-text file. We define a hit as the signal with the highest S/N per channel over all the drift rates searched. Only the signal with the highest S/N within a window  $\pm \dot{\nu}_{max} \times t_{obs}/2 = \pm 600$  Hz is recorded as a hit.

We used the Python PANDAS<sup>17</sup> package to read these files into a searchable table. To quickly process multiple files, we used DASK<sup>18</sup> to batch-process multiple files in parallel; this is far less computationally intensive than *turbo*SETI, and as such, only a single compute node was required. To load data and read observational metadata from filterbank and hdf5 formats, we use the BLIMPY<sup>19</sup> package (Price et al. 2019).

# 3.2. Data Selection

After *turbo*SETI is run on each file, sets of files are grouped to form complete ABABAB or ABACAD cadences. Observation sets that are not part of a complete cadence are not

<sup>&</sup>lt;sup>15</sup> https://github.com/ucberkeleyseti/turbo\_seti

<sup>&</sup>lt;sup>16</sup> https://github.com/ucberkeleyseti/tiddalik

<sup>&</sup>lt;sup>17</sup> https://pandas.pydata.org

<sup>&</sup>lt;sup>18</sup> https://dask.pydata.org

<sup>&</sup>lt;sup>19</sup> https://github.com/ucberkeleyseti/blimpy

<sup>&</sup>lt;sup>14</sup> https://breakthroughinitiatives.org/opendatasearch

analyzed further in this work. We require that all files with a cadence contain 16 integrations (5 minutes). We then use the find\_event method of *turboSETI* to search for hits that are present in all ON source observations, but not in OFF source observations.

We refer to hits matching this criterion as an "event." Specifically, any set of hits present in all ON observations in a frequency range calculated by  $2\dot{\nu}_0 \times t_{obs}$  and central frequency  $\nu_n$  are selected, where  $\dot{\nu}_{min} \leq |\dot{\nu}_0| \leq \dot{\nu}_{max}$  is the drift rate of the hit in the first ON observation, and  $\nu_n = \nu_0 + \dot{\nu}_0 \times (t_n - t_0)$ , while  $t_0$  and  $t_n$  are the observing start times of the first and *n*th observations, respectively. Additionally, any set of hits for which there is at least one hit in the OFF observations within  $\pm 600$  Hz of the hit frequency from the first ON observation would be discarded. This window corresponds to half the maximum searched drift of a signal over the period of the observation.

# 3.3. Event Grouping

Discrete analysis of an event without regard to surrounding events does not provide a complete picture. Events that are clustered in frequency, all of which exhibit the same drift rate, are likely to be associated with a single source of interference (or, indeed, technosignature).

We apply a simple grouping algorithm to assign events into groups, to aid in visualization and analysis. Events are grouped into frequency bins of width 125 kHz, then in each bin, the spacing between the highest and lowest start frequency is computed, to compute an effective bandwidth  $\Delta \nu_{\text{event}}$ , and central frequency  $\nu_{\text{event}}$ . We refer to each cluster as an "event group."

#### 3.4. Event Rejection and Analysis

We reject event groups where frequencies are within the GBT *L*-band or GBT *S*-band notch filters, but do not outright reject zero-drift signals. Examples of events that pass all criteria are shown in the Appendix. Any event group that satisfies these filters is considered as a candidate signal and visually inspected. To do so, we plot the dynamic spectra of the events using the BLIMPY package, for all on-source and reference pointings. We reject events when it is clear by eye that the event group is present in reference pointings, but was not detected above the *turbo*SETI S/N >10 threshold.

# 4. Results

We ran *turbo*SETI on all files with a complete observing cadence, finding a total of 51.71 million hits across the *L*-band, *S*-band, and 10 cm data sets (Table 4). Of these, 21,117 events were detected only in ON observations, which we clustered into 6154 event groups. We treated analysis for each receiver separately, detailed below.

### 4.1. Green Bank

*turbo*SETI was run on 1013 *L*-band observing cadences (506.5 hr) and 1076 *S*-band cadences (538.0 hr). We required that cadences had a full set of ON and OFF observations (ABACAD), that all subbands were present, and all had a cadence of 16 integrations, indicating that the observation was not interrupted.

 Table 4

 Summary of Hits (Signals above Threshold), Events (Hits Only in ON

 Observations), Event Groups (Clusters of Events), and Final Events (Groups with Limited Frequency Extent)

Receiver	Hits	Events	Groups	Final
GBT L band	37.14M	15998	4522	0
GBT S band	10.12M	5102	1572	0
Parkes 10 cm	4.45M	77	60	0

#### 4.1.1. L Band (1.10–1.90 GHz)

At the *L* band, a total of 37.14M hits were detected by *turbo*SETI above the S/N > 10 threshold. The hit distribution as a function of drift rate and S/N is shown in Figures 3(a) and (b), respectively. Dashed lines at  $\pm 0.25$  Hz s<sup>-1</sup> correspond to the maximum drift rate expected due to Earth's motion at the receiver's upper limit of 1.9 GHz. A majority of hits (21.90M) are at zero drift, a large fraction (13.9M) have negative drift rates, and a smaller fraction (1.37M) have positive drift rates. The bias of signals toward negative drift rates is likely due to satellites in non-geosynchronous orbits, which accelerate with respect to the telescope (Zhang et al. 2006).

Hit density as a function of frequency is shown in Figure 4(a). Areas of high hit density are associated with known sources of RFI, in particular Global Positioning Systems satellites (GPS), operating at 1.155–1.196 GHz (L5) and 1.555–1.596 GHz (L1). Bands associated with *GLONASS* satellites at 1.592–1.619 GHz (L1) and 1.192–1.212 GHz (L3) also show high hit density, along with the satellite downlink band 1.525–1.560 GHz.

Of the 37.14M hits, 15,998 events (hits found only in ON observations) were detected (Figures 3(a) and (b), light blue). This cut also excludes events within the GBT *L*-band notch filter (1.200–1.341 GHz). Events are concentrated at S/N < 1000, and at positive drift rates.

The events are clustered into 4522 event groups, from 831 unique stars. After visual inspection of these event groups, we do not find any signals that cannot be attributed to RFI; see the Appendix.

### 4.1.2. S Band (1.80-2.80 GHz)

At the S band, a total of 10.12M hits were detected; hit density as a function of drift rate and S/N is shown in Figures 3(c) and (d), respectively. A majority of hits (7.36M out of 10.1M) are at zero drift; however, a larger portion exhibit nonzero drift, with a small skew toward negative drift rates (1.55M versus 1.21M). Dashed lines at  $\pm 0.36$  correspond to the maximum drift rate expected due to Earth's motion at the receiver's upper limit of 2.7 GHz.

Of the 10.12M hits, 5102 events (hits found only in ON observations) were detected (Figures 3(c) and (d), light blue). Events are concentrated at low S/N (<10<sup>3</sup>) and at positive drift rates. The peaks at  $\pm 3.5$  Hz s<sup>-1</sup> apparent in Figure 3(c) are associated with RFI around 1930.2 MHz, within the cellular band, and as such may be due to cell phone activity near the observatory.

The hit density falls as S/N increases (Figure 3(d)), from  $\sim$ millions of hits per bin at S/N of 10, down to  $\sim$ tens of hits at S/N of 10<sup>7</sup>. Hit density as a function of frequency is shown in Figures 4(a) and (b). A large number of hits are attributable to RFI (see Table 6).

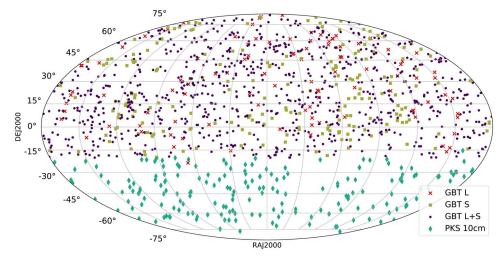


Figure 1. Distribution of observed sources in equatorial coordinates, taken from the 1702 star sample of Isaacson et al. (2017). Sources observed with Green Bank at both L band and S band are plotted in purple, sources only observed at the L band are plotted with red crosses, sources only observed at the S band are plotted with yellow squares, and sources observed with Parkes at 10 cm are plotted with aqua diamonds.

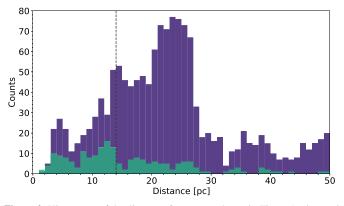


Figure 2. Histograms of the distances for sources shown in Figure 1, observed with Green Bank (purple) and Parkes (aqua).

After applying event grouping, a total of 5102 event groups were identified. After visual inspection, we do not find any candidate signals not attributable to RFI (Appendix).

#### 4.2. Parkes 10 cm (2.60–3.45 GHz)

We ran *turbo*SETI on the 966 observation cadences (483.0 hr); a total of 4.45M hits were detected. Histograms showing hit density as a function of drift rate are shown in Figures 3(e) and (f). A majority of hits (4.16M out of 4.45M) are at zero or within  $\pm 0.015$  Hz (Figure 3(e)). Outside of zero drift, there is a slight skew toward negative drift rates (134k versus 126k). Dashed lines at  $\pm 0.45$  correspond to the maximum drift rate expected due to Earth's motion at the 3.45 GHz upper limit of the receiver.

The hit density over S/N (Figure 3(f)) falls steadily until an S/N of  $\sim$ 1000, after which it rises rapidly before falling again. This may indicate distinct populations of interferers with different characteristic signal strengths that are not isotropically distributed.

Hit density as a function of frequency is shown in Figure 4(c). A large number of hits are associated with known RFI sources: 2.60–2.62 GHz is 4G cellular service downlink (band 7), and the 3.4 GHz band is licensed to the Australian National Broadband Network (NBN).

Of the 4.45M hits, only 60 event groups from 20 stars pass our selection criteria (Section 3.2). These events predominantly fall in the 3.40–3.45 GHz band and are likely interference associated with NBN. Example events are shown in Figure 9. In all cases, the narrowband signal was detected with an S/N > 10 in the ON source pointings, but not in the OFF source pointings. The effectiveness of the ON–OFF approach using Parkes at these frequencies indicates that this band is relatively quiet and, in regions where it is not, at least relatively stable.

### 5. Discussion

### 5.1. Comparison across Receivers

A summary of hits across the *L*-band, *S*-band, and 10 cm data sets is shown in Table 4. Broadly, the number of hits decreases with increasing receiver frequency. Taking into account receiver bandwidth and observation time, we can compare "hit rate density," that is, the number of hits per unit bandwidth per hour. This, and similarly defined event rate densities, is shown in Table 5.

A higher hit rate density corresponds to higher levels of RFI occupancy. However, hit rate is dependent on the S/N threshold, the sensitivity of the telescope, and observation strategy: a less sensitive telescope will report fewer hits above a given S/N for the same RFI environment; similarly, observations toward the horizon or known RFI sources will have higher rate densities. As such, direct comparison between observatories and observing campaigns is nuanced.

Nevertheless, the rate densities in Table 5 do inform us about broad RFI trends. The *L*-band receiver has a hit rate density more than  $5 \times$  higher than that of the *S*-band receiver, and  $10 \times$ that of the Parkes 10 cm receiver. Lower rate densities with the *S*-band receiver (for the same observational approach) indicate the RFI environment is cleaner over 1.9–2.8 GHz than 1.1–1.9 GHz at the GBT. The grouped event rate density is  $135 \times$  higher with the GBT *L* band than the Parkes 10 cm receiver, showing that a larger fraction of events pass our RFI identification at the *L* band than at 10 cm. We attribute this to a lack of satellites, or other rapidly moving sources, that transmit over 2.6–3.4 GHz.

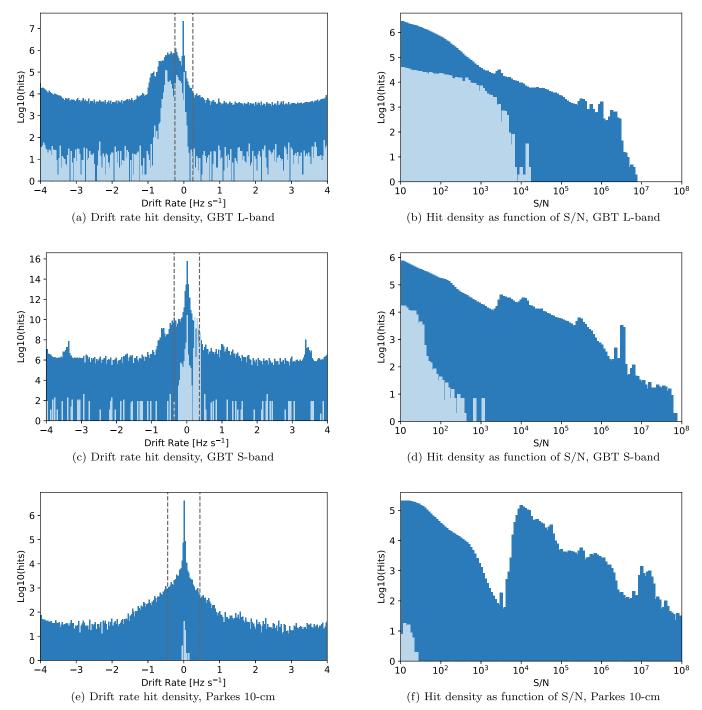


Figure 3. Histograms of hit (dark blue) and event (light blue) densities for GBT L-band, GBT S-band, and Parkes 10 cm observations, for drift rate (left) and as a function of signal to noise (right).

Plots of hit density as a function of frequency are given in Figure 4. Areas of high hit density are shaded; corresponding federal allocations are given in Table 6. In the US, the Federal Communications Commission (FCC<sup>20</sup>) and National Telecommunications and Information Administration (NTIA<sup>21</sup>) oversee spectrum allocations; in Australia, this is done by the Australian Communications and Media Authority (ACMA<sup>22</sup>). These

agencies coordinate with the International Telecommunications Union (ITU<sup>23</sup>), who also regulate space-based frequency allocations.

# 5.2. Combined Project Metrics and Figures of Merit

The search space for SETI signals is vast. Tarter (2001) describe the search space as a "nine-dimensional haystack"; this metaphor is continued in Wright et al. (2018), who detail a method to compute a "haystack fraction" that quantifies how

<sup>&</sup>lt;sup>20</sup> http://fcc.gov

<sup>&</sup>lt;sup>21</sup> http://ntia.gov

<sup>&</sup>lt;sup>22</sup> http://acma.gov.au

<sup>&</sup>lt;sup>23</sup> https://www.itu.int

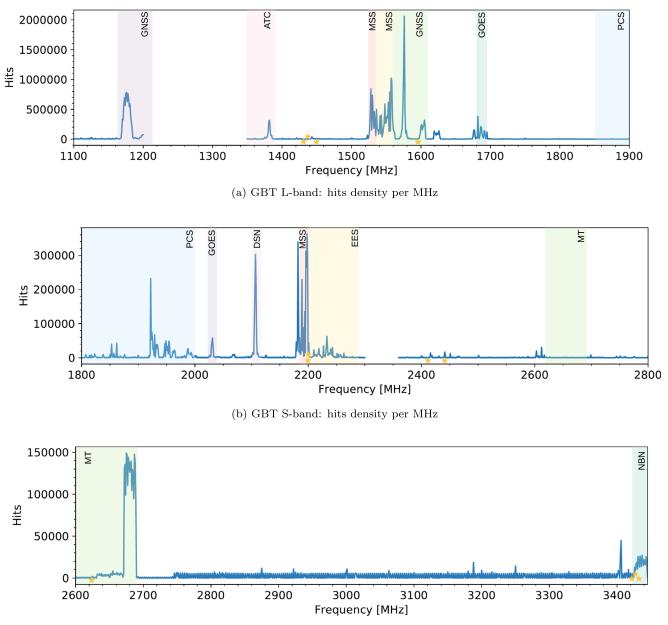




Figure 4. Frequency distribution of hits produced by the *turbo*SETI pipeline, over the bands of the GBT *L*-band, GBT *S*-band, and Parkes 10 cm receivers. Frequencies with large hit densities (shaded) are associated with known sources of RFI; see Table 6. Events of interest as detailed in the Appendix are marked with stars.

Table 5
Rate Densities for Hits (Signals above Threshold), Events (Hits Only in ON
Observations), and Event Groups (Clusters of Events)

	L Band	S Band	10 cm
Hit Rate Density ( $\times 10^{-6}$ Hits Hz <sup>-1</sup> hr <sup>-1</sup> )	111.1	20.0	10.8
Event Rate Density $(\times 10^{-9} \text{ Events Hz}^{-1} \text{ hr}^{-1})$	47.9	10.1	0.18
Grouped Event Rate Density $(\times 10^{-9} \text{ Groups Hz}^{-1}\text{hr}^{-1})$	13.5	3.1	0.1

complete a SETI search is. The haystack fraction is but one of several Figures of Merit (FoM) that can be used as heuristics to compare SETI searches.

An historical FoM is the Drake FoM (DFM; Drake 1984), which is given by

$$\text{DFM} = \frac{\Delta \nu_{\text{tot}} \Omega}{F_{\text{min}}^{3/2}},\tag{4}$$

Price et al.

where  $\Delta \nu_{\text{tot}}$  is the observing bandwidth,  $\Omega$  is the sky coverage, and  $F_{\text{min}}$  is the minimum detectable flux in W m<sup>-2</sup>. The -3/2 index on  $F_{\text{min}}$  encompasses distance-to-volume scaling ( $d^3$ ) and sensitivity scaling ( $d^{-2}$ ).

As we use three receivers with varying fields of view, system temperatures, and bandwidths, we compute a combined  $DFM_{tot}$ :

$$DFM_{tot} = \sum_{i}^{N} DFM_{i},$$
(5)

 Table 6

 Radio Frequency Spectrum Allocations, for Bands with High Hit Densities (see Figure 4)

Band (MHz)	Federal Allocation
1164–1215	Global Navigation Satellite System (GNSS) <sup>a</sup>
1350-1390	Air traffic control (ATC) <sup>a</sup>
1525-1535	Mobile-satellite service (MSS) <sup>a</sup>
1535-1559	MSS <sup>a</sup>
1559-1610	GNSS <sup>a</sup>
1675–1695	Geostationary operational environmental satellite (GOES) <sup>a</sup>
1850-2000	Personal communications services (PCS) <sup>a</sup>
2025-2035	GOES <sup>a</sup>
2100-2120	NASA Deep Space Network use (DSN) <sup>a</sup>
2180-2200	MSS <sup>a</sup>
2200-2290	Earth exploration satellite (EES) <sup>a</sup>
2620-2690	Mobile Telecommunications (MT) <sup>b</sup>
3425-3492.5	National Broadband Network (NBN) <sup>b</sup>

Notes.

<sup>a</sup> FCC/NTIA Table of Frequency Allocations.

<sup>b</sup> ACMA Australian Radiofrequency Spectrum Plan.

i.e., the sum of DFMs; larger DFM values are better. The DFM<sub>tot</sub> for this project is  $9.2 \times$  than that of Enriquez et al. (2017). The combined sky coverage for all the observations was 22.1 deg<sup>2</sup>, in contrast with the 10.6 deg<sup>2</sup> presented by Enriquez et al. (2017).

We note that the above formulation of the DFM assumes a common channel bandwidth. For a narrowband signal,  $F_{min}$  depends upon the signal-to-noise threshold S/N<sub>min</sub> and may be calculated as

$$F_{\rm min} = S/N_{\rm min} \frac{2k_{\rm B}T_{\rm sys}}{A_{\rm eff}} \sqrt{\frac{B}{n_{\rm pol}t_{\rm obs}}},$$
(6)

where  $k_{\rm B}$  is the Boltzmann constant,  $T_{\rm sys}$  is the system temperature,  $A_{\rm eff}$  is the effective collecting area of the telescope, *B* is the channel bandwidth, and  $n_{\rm pol}$  is the number of polarizations. Note that  $F_{\rm min}$  (flux) is related to flux density  $F_{\rm min} = S_{\rm min} / \delta \nu_t$ , where  $\delta \nu_t$  is the bandwidth of the transmitting signal. We have chosen a value of unity in this work.

Wright et al. (2018) presents a formalism in which one defines "boundaries" to specify an *N*-dimensional survey space, or "haystack." One can then compute what fraction of a given haystack a survey probes. Using their boundaries, Wright et al. (2018) compute a haystack fraction of  $3.8 \times 10^{-19}$  for the observations presented in Enriquez et al. (2017). For our *L*-band, *S*-band, and 10 cm observations, we compute haystack fractions  $1.23 \times 10^{-18}$ ,  $7.44 \times 10^{-19}$ , and  $3.37 \times 10^{-19}$ . These correspond to  $3.24 \times$ ,  $1.96 \times$ , and  $0.88 \times$  times those of Enriquez et al. (2017), respectively. The integrated value taking all the observations together is  $6.08 \times$ .

The DFM and haystack fraction are useful heuristics when comparing surveys. However, neither the DFM nor the haystack fraction take into account the distance to survey targets: they treat observations of nearby stars equal to a patch of seemingly blank sky. For this reason, Enriquez et al. (2017) define the "Continuous Waveform Transmitter Rate Figure of Merit," or TFM:

$$\text{TFM} = \eta \frac{\text{EIRP}_{\text{min}}}{N_{\text{star}}} \frac{\nu_c}{\Delta \nu_{\text{tot}}},\tag{7}$$

where  $\nu_c$  is the central observing frequency,  $N_{\text{star}}$  is the number of stars observed, EIRP<sub>min</sub> is the minimum detectable equivalent isotropic radiated power (EIRP, in W), and  $\eta$  is a normalization factor. The value  $(\nu_{\text{tot}}/\nu_c)/N_{\text{star}}$  encompasses fractional bandwidth and number of sources, and is referred to as the transmitter rate. The EIRP<sub>min</sub> for a given target increases with the distance squared:

$$\mathrm{EIRP}_{\mathrm{min}} = 4\pi d^2 F_{\mathrm{min}}.$$
 (8)

The EIRP<sub>min</sub> for our GBT observations is  $2.1 \times 10^{12}$  W, and  $9.1 \times 10^{12}$  W for Parkes observations at the 50 pc maximum distance of the Isaacson et al. (2017) sample.

Numerically lower TFM scores represent more sensitive and more complete surveys. For comparison, the TFM for this work is  $6.95 \times$  smaller than that for Enriquez et al. (2017); comparisons against other surveys are shown in Figure 5.

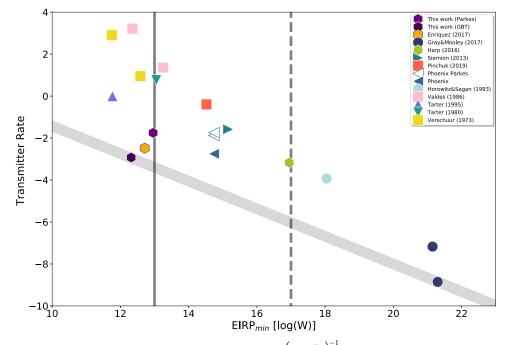
### 5.3. Limits on Narrowband Technosignatures

We find no evidence for narrowband transmitters from observations of our target stars above the EIRP<sub>min</sub> values of  $2.1 \times 10^{12}$  W for GBT observations and  $9.1 \times 10^{12}$  W for Parkes observations.

It is difficult to place limits on the existence of putative transmitters in the direction of the star targets, due to the presence of RFI, potential intermittency/periodicity of the transmission, or our data analysis being insensitive to a given signal due to pipeline limitations (see Section 5.7). Nevertheless, one may still place a probabilistic upper limit on the prevalence of putative continuous narrowband transmitters above EIRP<sub>min</sub>, assuming that such transmitters are rare. That is, one may compute a conditional probability of detecting a signal, should it exist above EIRP<sub>min</sub> and within the observing band, by treating each star target as a trial within a Poissonian distribution.

We make a conservative estimate that a given observation has a probability of P = 0.5—to account for potential RFI obscuration—of detecting a narrowband signal at a random frequency within the observed band (above the EIRP<sub>min</sub>). For the GBT *L* band, 882 star targets were observed; treating these as discrete trials, we place a limit with 95% confidence that fewer than 0.45% of stars have narrowband transmitters above an EIRP<sub>min</sub> of  $2.1 \times 10^{12}$  W. For the *S* band, this limit is 0.37% of stars (EIRP<sub>min</sub> $2.1 \times 10^{12}$  W) and 2.0% based on the Parkes 10 cm data (EIRP<sub>min</sub> $9.1 \times 10^{12}$  W).

These limits are coarse and could be improved by more careful consideration of several aspects. First, one could inject signals at various drift rates into real observational data to compute signal recovery statistics. Second, one could run Monte Carlo simulations in which transmitters are placed at different distances, and their signal properties are drawn from varied probability distributions, following a Bayesian approach (e.g., Grimaldi & Marcy 2018); this is an avenue for future investigation.



**Figure 5.** Transmitter rate comparison with other historical projects. The transmitter rate,  $\left(N_{\text{star}}\left(\frac{\nu_c}{\nu_{\text{tot}}}\right)^{-1}\right)$ , is plotted on logarithmic axes against the minimum detectable power, EIRP<sub>min</sub>, based on the distance to the farthest star in the sample. Points toward the bottom of this plot represent surveys with large numbers of star targets and high fractional bandwidth; points toward the left represent surveys where the product of sensitivity and distance to targets is lower. The solid and dashed vertical lines represent the EIRP of the Arecibo planetary radar, and the total power from the Sun incident on Earth, respectively. The diagonal gray line is a fit between the most constraining data points for the transmitter rate and EIRP<sub>min</sub>.

### 5.4. Comparison to Enriquez et. al. (2017)

In addition to new observations, we reanalyzed the observations reported by Enriquez et al. (2017) over 1.10–1.90 GHz, using a lower S/N cutoff (10 versus 25). The result is that our sensitivity is better by a factor of  $2.5 \times$ , but as a side effect, our false-positive rate also increases. A number of events were recorded where a signal was present in the OFF observation, but was not detected above the S/N threshold; Enriquez et al. (2017) avoided this by requiring a higher S/N (25) for ON observations than OFF observations (20). We also reanalyzed data across a broader range of drift rates, expanding from  $\pm 2$  to  $\pm 4$  Hz s<sup>-1</sup>. A side effect of the larger drift rate is that the window used to avoid redundant detections, given by  $\pm \dot{\nu}_{max} \times t_{obs}/2$ , is larger.

### 5.5. Comparison to Previous 1.80–3.45 GHz Searches

SETI searches over the combined 1.80-3.45 GHz range of the GBT S-band and Parkes 10 cm receiver have been conducted previously, but to a lesser extent than the so-called "water hole" between 1.42 and 1.67 GHz. As part of the SERENDIP-II survey, Werthimer et al. (1986) observed a 32 kHz band around 2.25 GHz using a 210 ft antenna. Project Sentinel (Horowitz & Forster 1985) and project META (Horowitz & Sagan 1993) ran narrowband searches around 2.84 GHz (twice the 21 cm line frequency). Tarter & Klein (1995) observed 24 solar-type stars using a system with 8 MHz bandwidth at 8 cm and 12 cm wavelengths as part of the High Resolution Microwave Survey; this survey was defunded by congress before completion. Project Phoenix also covered this band, but published details on the observations are sparse (Backus & Project Phoenix Team 2004). To our knowledge, no archival data from any of these projects are publicly available.

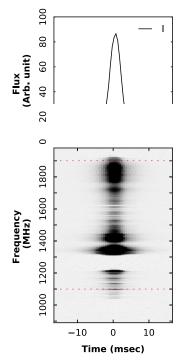
More recently, the Allen Telescope Array (Harp et al. 2016) observed 9293 stars sporadically over 1–9 GHz (averaging 785 MHz of observed band per star) and report a minimum detectable flux density  $S_{\rm min} = 271$  Jy at 3 GHz. This is roughly 7× more primary targets, but at  $17 \times -39 \times$  lower sensitivity, with a lower range of drift rates searched. By Equation (6), to reach an equivalent sensitivity would require observations between 289 and 1521 times longer. Harp et al. (2016) searched drift rates of  $\pm 2$  Hz s<sup>-1</sup>, over the 1.1–3.4 GHz band.

### 5.6. Comparison to Other Recent GBT Searches

Recent SETI searches were also undertaken at the GBT by Margot et al. (2018) and in follow-up work by Pinchuk et al. (2019; henceforth M&P). In both cases, the GBT *L*-band receiver was used, but different data analysis approaches were applied. Careful comparison of the two approaches to identify their relative advantages is invaluable. Here, we discuss the differences and similarities between these campaigns and the BL *L*-band observations to identify areas where future analyses can be improved.

### 5.6.1. Observational Strategy

Both M&P and BL employed a similar observational strategy whereby targets were observed multiple times. M&P observed each target twice for 150 s per pointing, whereas we observed each target three times, for 300 s per pointing. Our target selection, detailed by Isaacson et al. (2017), draws from a morphologically diverse selection of stars, containing most types of stars existing within 50 pc; M&P selected targets with known exoplanets, predominantly GKM type stars from the *Kepler* field, as well as two nearby planet-hosting M dwarfs.



**Figure 6.** Observation of PSR J0826+2637 at the *L* band from the GBT using the BL backend. The top panel shows the integrated profile while the bottom panel shows dynamic spectra. The notch filter is clearly evident at frequencies between 1.2 and 1.33 GHz. The red dotted lines show the edges of the frequency band between 1.1 and 1.9 GHz. A very clear pulsar detection can be seen well across these limits.

Combined, a total of 26 targets were observed in the M&P sample, over 130 minutes.

Margot et al. (2018) and Pinchuk et al. (2019) only analyzed data within the nominal 1.15–1.73 GHz passband of the receiver. Apart from elevated system temperature due to loss in aperture efficiency, we find no impediment to the use of the full 1.10–1.90 GHz band, although we note that both BL and M&P avail themselves of the 1.20–1.34 GHz notch filter to suppress nearby air surveillance radar (see Figure 6).

### 5.6.2. 2 Bit Requantization

The most significant differences arise at the data recording level. M&P employs the older GUPPI processing system, which records data as 2 bit quantized voltages (Siemion et al. 2013; Prestage et al. 2015). In contrast, BL records data at 8 bit resolution. BL then converts the recorded voltages into spectral products, resulting in a  $\sim$ 50× reduction in data volume. While BL can archive 8 bit voltages, only a subset of voltages are retained due to storage limitations. Also, while voltage-level data products are more flexible than (fixed-resolution) spectral products, storing 2 bit data would require a 12.5× increase in storage capacity (or 12.5× decrease in observing time).

Requantizing to lower bit depth has several negative effects. The first is that the dynamic range—the ratio between largest and smallest possible values—is limited. The dynamic range in decibels available within N bits is<sup>24</sup>

$$DR = 20 \log_{10}(2^N - 1) (dB).$$
(9)

For 2 and 8 bit data, the dynamic range afforded is 9.54 dB and 48.13 dB, respectively. Any signal that saturates the available dynamic range will be distorted. Impulsive RFI can introduce harmonic distortions and interfere with nominal requantization. During requantization, the system bandpass must generally be subtracted dynamically, using scaling factors that change over small (~seconds) time windows. If the scaling factors are not preserved, bandpass information—a useful diagnostic—is lost.

Quantization efficiency—the relative loss in S/N due to quantization—drops from 0.99912 for 8 bit down to 0.881154 for 2 bit data, assuming optimal level settings (Thompson et al. 2007). From a SETI perspective, the end result of requantization to 2 bits is that S/N threshold (e.g., S/N > 10) would need to be lowered (S/N > 8.8) to retrieve the same number of hits. The limited dynamic range also places a limit on the maximum S/N achievable, and strong hits will exhibit harmonic distortions that may register as extraneous hits. For these reasons, in situations where we store voltage-level data, we retain an 8 bit resolution.

### 5.6.3. Drift Rates and S/N Threshold

Margot et al. (2018) use a window of size  $\pm \dot{\nu}_{max} \times t_{obs} \approx \pm 1500$  Hz to remove redundant detections. Following Enriquez et al. (2017), we use  $\pm \dot{\nu}_{max} \times t_{obs}/2$ , i.e., half this value, but we fix the window size to  $\pm 600$  Hz when comparing ON and OFF observations for RFI rejection. An issue with such maximum-drift-based approaches is that as the drift rate searched increases, the fraction of "blanked" band also increases, so candidate signals may be discarded, and metrics such as the DFM may be overestimated. To combat this, Pinchuk et al. (2019) instead require that detections do not cross in time–frequency space. For continuity with Enriquez et al. (2017), we do not implement such a strategy here.

Pinchuk et al. (2019) estimate that the Margot et al. (2018) DFM was overestimated by  $\sim 5\%$  due to blanking. We compute average "blanked" fractions of 0.9%, 0.2%, and 0.1% for the GBT *L*, GBT *S*, and Parkes 10 cm observations, excluding notch filters and a 1 kHz region around each hit, so this effect is negligible. Signal rejection filters will also affect metrics such as the DFM; we emphasize that the FoM should be treated as heuristics for comparison of observational campaigns only.

### 5.6.4. Event Grouping and Rejection

Pinchuk et al. (2019) grouped hits into  $\sim$ kHz bins and discarded all hits in bins with high hit density. This is similar to our event grouping approach; however, we promote events for visual inspection. Both of these approaches could likely be improved by identifying other signal properties (e.g., bandwidth, kurtosis) and forming a larger-dimensional parameter space in which to cluster signals.

### 5.7. Pipeline Limitations

Based on the analysis of events, we identify several limitations of our current pipeline and corresponding areas of improvement. First, it is often clear by visual analysis that a signal is indeed present in an OFF observation but was below the S/N threshold required. The false-positive rate can be decreased by setting variable thresholds for OFF source pointings that account for the fact that sources well off axis to the observing direction can nevertheless have varying

 $<sup>^{24}</sup>$  https://www.analog.com/media/en/training-seminars/tutorials/MT-229.pdf

apparent power in our ON and OFF source positions. Another possible method is to compute the cross-correlation between ON and OFF pointings over a set of lags and search directly for signals present common in ON and OFF at a lower total threshold.

Due to how the S/N is calculated, the S/N for events with bandwidth greater than a single channel is underestimated. By decimating in frequency (i.e., averaging over steps of  $2^N$  channels), the S/N for wider bandwidth signals will increase until the signal is no longer resolved in frequency (e.g., as employed by Siemion et al. 2013). This approach is already used commonly in RFI flagging codes (e.g., Offringa et al. 2012). The estimation of S/N is also sensitive to the estimate of noise levels: in areas of high RFI occupancy, noise level estimates will be affected by the presence of RFI.

For frequency-resolved signals, the S/N can also be improved by averaging across the bandwidth of the signal. One could use a hierarchical frequency decimation approach, searching optimal drift rates ranges at each stage to ensure the drift rate does not exceed  $B/t_{obs}$ , which leads to smearing across channels. In a recent paper, Sheikh et al. (2019) advocates a rate as high as 200 nHz, so as to be sensitive to a larger class of bodies, including exoplanets with highly eccentric orbits (e.g., HD 80606b) and small semimajor axes such as Kepler-78 b.

Our frequency and drift rate grouping algorithm is simplistic and could be improved using methods from machine learning (ML), such as *k*-means clustering. Drift rate and frequency are only two signal properties that could be used for grouping events, and with proper labeling, grouping could also take into account bandwidth, signal kurtosis, or other assessments of modulation type. With appropriate training, ML methods can also be used to self-identify features (Zhang et al. 2019).

## 6. Conclusions

As part of the BL program, we searched 1327 nearby stars taken from the Isaacson et al. (2017) sample for technosignatures, using data from the Green Bank and Parkes telescopes. We used three receivers, spanning a combined range of 1.10–3.45 GHz, and found no compelling candidates that are not attributable to RFI. Our *turbo*SETI pipeline searched for narrowband signals exhibiting time-variable frequency drift due to Doppler acceleration, finding over 51 million hits above our S/N threshold. Of these hits, we identified 6154 event groups that passed our automated verification tests; however, none of these passed closer manual inspection and cross-referencing against known RFI.

Combined, these observations constitute the most comprehensive survey for radio evidence of advanced life around nearby stars ever undertaken, improving on the results of Enriquez et al. (2017) in both sensitivity and number of stars. Together with other recent work from the resurgent SETI community, we are beginning to put rigorous and clearly defined limits on the behavior of advanced life in the universe. We note that significant additional observational and theoretical work remains to be done before we are able to make general statements about the prevalence of technologically capable species.

With respect to the specific search described here, our analysis is currently confined to only spectrally narrow drifting signals using our highest resolution data product. A high-timeresolution data product will be searched for pulsed signals in future work, and a refined drifting spectral line search will be undertaken covering wider bandwidths. Further, BL observations with the GBT and Parkes are ongoing, with the GBT *C*band (3.9–8.0 GHz), GBT *X*-band (8.0–11.6 GHz), and Parkes UWL (0.7–4.0 GHz) receivers. Observations of the Galactic Plane are also being undertaken, using the Parkes multibeam receiver (1.2–1.6 GHz). Observations are also planned with the MeerKAT telescope and other partner facilities such as the Murchison Widefield Array.

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*Software:* BLIMPY (Price et al. 2019), *turboSETI* (Enriquez & Price 2019), ASTROPY (Price-Whelan et al. 2018), H5PY (Collette 2013), DASK (Dask Development Team 2016), PANDAS (McKinney et al. 2010), MATPLOTLIB (Hunter 2007), NUMPY (Oliphant 2006), SCIPY (Jones et al. 2001), JUPYTER (Kluyver et al. 2016).

# Appendix Example Events

A number of events passed our automated verification tests but failed manual inspection. In this appendix, we discuss these events in further detail. Figures 7–9 show the four top-ranked events from each of the *L*-band, *S*-band, and 10 cm data sets. These events are marked in Figure 4 with gold stars.

At the L band (Figure 7)—Two of the most compelling events were from observations of HIP 54677 (Figures 7(a) and (b)), appearing only in ON observations. These two signals are spaced  $\sim$ 3 MHz apart, have the same drift rate, similar bandwidth, and similar power levels. A similar event to Figure 7(a) was detected  $\sim$ 1 kHz lower in an observation of HIP 103388; however, this event was classified as RFI due to the presence of signal in OFF observations. In addition, RFI events were detected within 40 Hz of the central frequency in Figure 7(b). Taken together, these RFI events indicate that the HIP 54677 events are also RFI. These events fall in the 1435–1525 MHz band used for aeronautical telemetry.<sup>25</sup>

As shown in Figure 7(c), *turbo*SETI detected an event only visible in ON pointings toward HIP 100064. A second drifting signal can be seen within the plotted band; however, this is detected in both ON and OFF pointings. We reject this event as a similar pair of drifting signals are seen in observations of HIP 32423, within 30 Hz of the central frequency of Figure 7(c); a similar pair is also seen in observations of HIP 21402,  $\sim$ 1.5 kHz above that of HIP 100064.

Figure 7(d) shows an event in the direction of HIP 1444. Events with similar signal bandwidths were detected  $\sim$ 2.4 kHz below (observations of HIP 56802) and  $\sim$ 1.6 kHz above (HIP 99572). However, these RFI events have differing drift rates and do not exhibit the change of drift rate present in the HIP 1444 event.

<sup>&</sup>lt;sup>25</sup> https://www.ntia.doc.gov/files/ntia/publications/.../1435.00-1525.00\_ 01DEC15.pdf

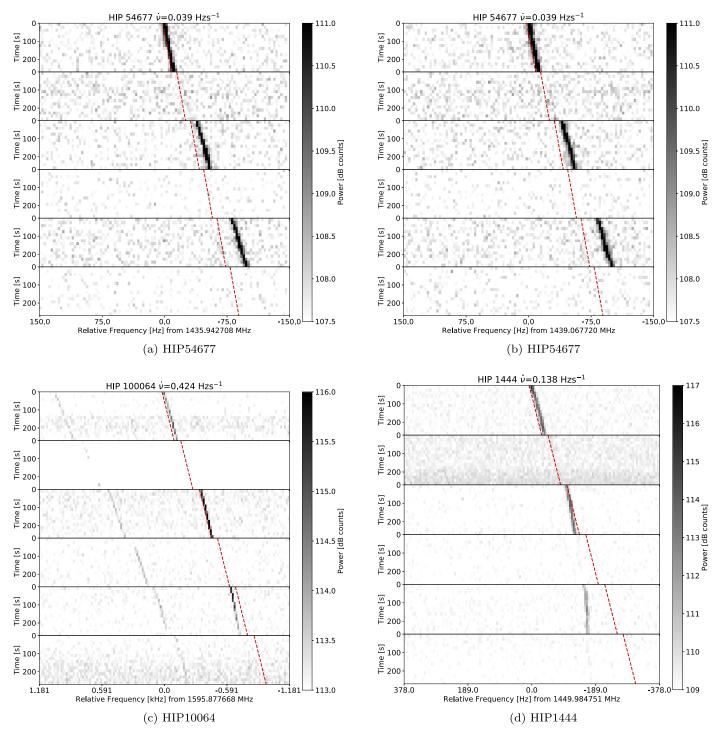


Figure 7. Dynamic spectrum for selected *turbo*SETI events from GBT *L*-band observations. Each subfigure shows a full ABACAD cadence; each of the six panels represents ON and OFF source, consecutively. The red dashed lines show the drift rate as detected by the pipeline; a small frequency offset has been added for visualization.

At the S band (Figure 8)—Narrowband signals with nonzero drift were detected above the S/N threshold in observations of HIP 91699 (Figure 8(a)) and HIP 22845 (Figure 8(b)). These events have durations under 5 minutes, so appear to turn on and off during the observations. A total of 22 other events with similar drift rates and frequency–time structure were found within 2200.04– 2200.5 MHz, some of which appear in OFF pointings; as such, we identify the HIP 91699 and HIP 22845 events as RFI. The 2200–2290 MHz band is used for spacecraft tracking and telemetry.<sup>26</sup>

The HIP 44072 event (Figure 8(c) displays complex structure in both time and frequency. Similar events were found in observations of HIP 68030, HIP 13402, and HIP 113178; as such, we identify this as RFI. Similarly,

<sup>&</sup>lt;sup>26</sup> https://www.ntia.doc.gov/files/ntia/publications/.../2200.00-2290.00-01MAR14.pdf

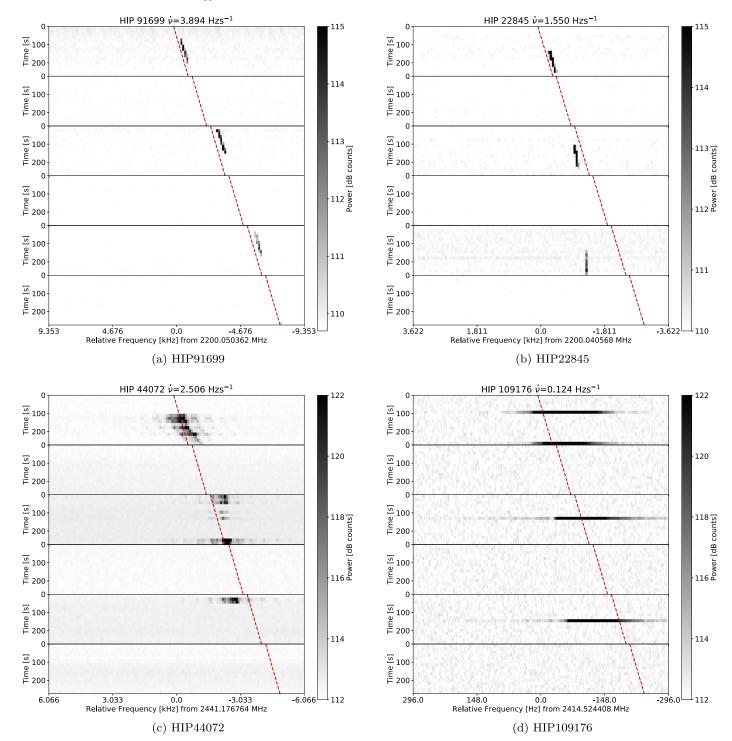


Figure 8. Dynamic spectrum for selected *turboSETI* events from GBT S-band observations. Each subfigure shows a full ABACAD cadence; each of the six panels represents ON and OFF source, consecutively. The red dashed lines show the drift rate as detected by the pipeline; a small frequency offset has been added for visualization.

HIP 109716 is similar to events detected in HIP 77655 pointings in which hits are visible in both ON and OFF pointings.

Parkes 10 cm receiver (Figure 9)—Remarkably fewer events were detected at Parkes, with only 60 event groups passing

automated verification. All of these signals are present in both ON and OFF observations, but were not detected above the required S/N of 10 in OFF observations. Given their frequency extent, most are likely associated with the NBN that is known to operate at 3.4 GHz.

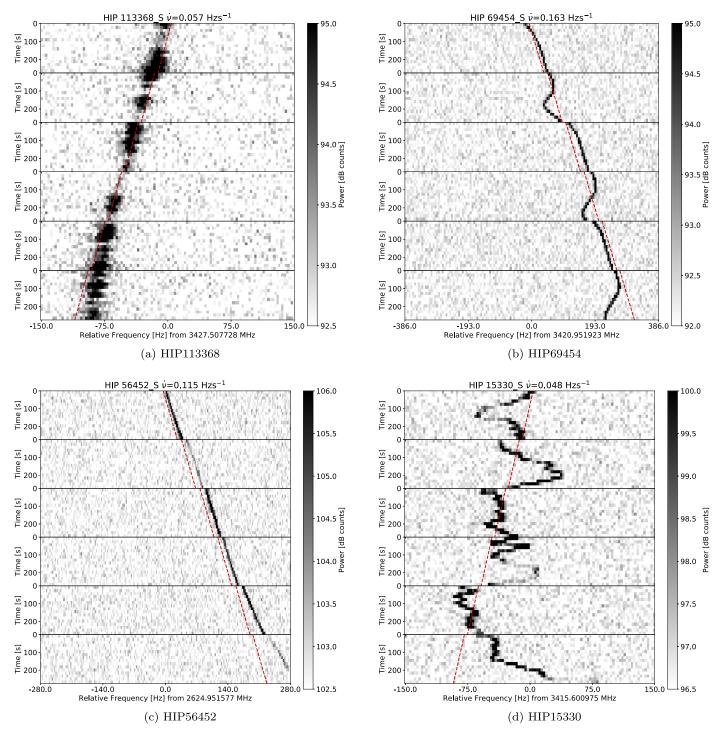


Figure 9. Dynamic spectrum for selected *turbo*SETI events from Parkes 10 cm observations. Each subfigure shows a full ABABAB cadence; each of the six panels represents ON and OFF source, consecutively. The red dashed lines show the drift rate as detected by the pipeline; a small frequency offset has been added for visualization.

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